

HNNLO version 2.0

This is a note about the HNNLO program. HNNLO is a parton level Monte Carlo program that computes the cross section for SM Higgs boson production in pp and $p\bar{p}$ collisions. The calculation is performed up to NNLO in QCD perturbation theory. The present version includes the decay modes $H \rightarrow \gamma\gamma$, $H \rightarrow WW \rightarrow l\nu l\nu$ and $H \rightarrow ZZ \rightarrow 4$ leptons. The user is allowed to apply arbitrary cuts on the final state and to plot the corresponding distributions in the form of bin histograms. If you use this program, please quote Refs. [1, 2, 3].

1 Introduction

The HNNLO program is based on an extension of the subtraction formalism to NNLO, as described in Ref. [1].

The calculation is organized in two parts. In the first part (virtual), the contribution of the regularized virtual corrections (up to two-loop order) is computed. In the second part (real), the cross section for the production of the Higgs boson in association with one jet is first evaluated up to NLO (i.e. up to $\mathcal{O}(\alpha_S^4)$). This step of the calculation can be performed with any available version of the subtraction method. Here we use the dipole formalism [4], as implemented in the MCFM Monte Carlo program [5]. Since the $H + \text{jet}$ cross section is divergent when the transverse momentum q_T of the Higgs boson becomes small, a suitable counterterm must be subtracted to make the result finite as $q_T \rightarrow 0$. The program uses the counterterm introduced in Ref. [1], and thus it completes the evaluation of the real part. Finally, the two contributions (virtual and real) can be combined to reconstruct the full cross section.

The calculation can be performed by using the large- M_t approximation, M_t being the mass of the top quark, or by keeping the exact dependence of the top- and bottom-quark masses up to NLO [3]. NNLO effects are always evaluated in the large- M_t limit, by rescaling the corresponding $\mathcal{O}(\alpha_S^4)$ contribution with the exact top-mass dependence at Born level.

The present version of the program includes the decay modes $H \rightarrow \gamma\gamma$ [1], $H \rightarrow WW \rightarrow l\nu l\nu$ and $H \rightarrow ZZ \rightarrow 4$ leptons [2]. In the latter case the user can choose between $H \rightarrow ZZ \rightarrow \mu^+\mu^-e^+e^-$ and $H \rightarrow ZZ \rightarrow e^+e^-e^+e^-$, which includes the appropriate interference contribution.

The program treats the Higgs boson in the narrow width approximation. In the decay modes $H \rightarrow WW \rightarrow l\nu l\nu$ and $H \rightarrow ZZ \rightarrow 4$ leptons, the finite width of the W and Z bosons is instead taken into account.

The program can be downloaded from <http://theory.fi.infn.it/grazzini>. To extract it, simply use `tar -xzvf hnnlo-v2.0.tgz` and the `hnnlo-v2.0` directory will be created. The structure of the directory is

- **bin**: The directory containing the executable `hnnlo` and the input and output files.
- **doc**: The directory containing this note.
- **obj**: The directory containing the object files.
- **src**: The directory containing the source of the code.

In the present version of the code we offer the possibility to use the program with the Les Houches Accord PDF (LHAPDF) interface. To do so, the user should set the variable `PDFROUTINES=LHAPDF` in the makefile and `LHAPDFLIB` to the path of the LHAPDF library.

2 Implementation of cuts

Before compiling the program, the user must choose the cuts to apply on the final state. This is done through the `cuts` and `isolation` subroutines. The default version of the subroutine `cuts` contains selection cuts that are used in the search for the Higgs boson at the LHC. In the default version, the lines of the code implementing the various cuts are all commented, thus the program will return the total cross section for the selected decay channel. Since the computation is performed in the narrow-width approximation, this will correspond to the total cross section $gg \rightarrow H$ times the branching ratio in the corresponding channel.

In order to activate the various cuts the user must uncomment the corresponding lines in `cuts.f`. Photon or lepton isolation can be implemented by switching to `true` the logical variable `isol`. The parameters to define the isolation procedure are set in the `isolation` subroutines.

Beside the leptons (photons) the program makes possible to identify the final state jets. Jets are reconstructed according to the k_T algorithm with $R = 0.4$. Different values of R can be implemented by modifying the `setup` subroutine¹.

Both `cuts.f` and `isolation.f` can be found in the `/src/User` directory. The `setup.f` source file can be found in the `/src/Need` directory.

3 Compilation

The program is self-contained and it has been successfully tested on Linux and Mac-OS X environments. To compile the code descend in the `hnnlo` directory and simply type

- `make`

To run it go in the `bin` directory and type:

- `hnnlo < infile`

¹Up to NNLO the k_T and anti- k_T algorithms are equivalent in this case.

4 The input file

This is a typical example of input file:

```
8d3 ! sroot
1 1 !ih1 ih2
125d0 ! hmass
125d0 125d0 ! mur, muf
2 ! order
1 ! higgsdec
'tota' ! part
15 8000000 ! itmx1, ncall1
30 8000000 ! itmx2, ncall2
617 ! rseed
92 0 ! iset, nset
172.5d0 ! mtop
4.75d0 ! mbot
2 ! approxim
'MSTW2008nnlo68cl.LHgrid' 0 ! set, member (LHAPDFs)
'nnlo' ! runstring
```

- **sroot**: Double precision variable for CM energy (GeV).
- **ih1, ih2**: Integers identifying the beam (proton=1, antiproton=-1)
- **hmass**: Higgs boson mass (GeV). This is a double precision variable that sets the mass of the SM Higgs boson;
- **mur, muf**: Renormalization and factorization scales (GeV): can be different from each other but always of order M_H .
- **order**: Integer setting the order of the calculation: LO (0), NLO (1), NNLO (2).
- **higgsdec**: Decay mode of the Higgs: $H \rightarrow \gamma\gamma$ (1), $H \rightarrow WW \rightarrow l\nu l\nu$ (2), $H \rightarrow ZZ \rightarrow e^+e^-\mu^+\mu^-$ (31), $H \rightarrow ZZ \rightarrow e^+e^-e^+e^-$ (32).
- **part**: String identifying the part of the calculation to be performed: **virt** for virtual contribution, **real** for real contribution, **tota** for the complete calculation.
- **itmx1, ncall1**: Number of iterations and calls to vegas for setting the grid.
- **itmx2, ncall2**: Number of iterations and calls to vegas for the main run.
- **rseed**: Random number seed.
- **iset, nset**: Integers identifying the pdf set chosen and the eigenvector for computing PDF errors (if the native PDF interface is used). A list of available pdf is given below.
- **mtop**: Mass of the top-quark (GeV)

- **mbot**: Mass of the bottom-quark (GeV)
- **approxim**: Integer defining the approximation to be used: large- M_t approximation (0), exact dependence on M_t up to NLO (1), exact dependence on M_t and M_b up to NLO (2).
- **PDFname**, **PDFmember**: String identifying the PDF set chosen and integer identifying the member for PDF errors (if the LHAPDF interface is used).
- **runstring**: String for grid and output files.

5 Output

At the end of the run, the program returns the cross section and its error. The program also writes an output file in the topdrawer format containing the required histograms with an estimate of the corresponding statistical errors. During the run, the user can control the intermediate results. The plots are defined in the `plotter` subroutine. The user can easily modify this subroutine according to his/her needs.

To obtain a cross section accurate at the percent level at NLO typically requires about one hour of run on a standard PC. At NNLO the required run time is at least about a factor of 10 larger.

6 Parton distributions

The HNNLO program can be compiled with its own Parton Distribution Functions (PDF) interface (set `PDFROUTINES = NATIVE` in the `makefile`) or with the LHAPDF interface (set `PDFROUTINES = LHAPDF` in the `makefile`). We point out that the value of $\alpha_S(m_Z)$ is not adjustable; it is hard-wired with the value of $\alpha_S(m_Z)$ in the parton distributions. Moreover, the choice of the parton distributions also specifies the number of loops that should be used in the running of α_S . A list of available parton densities for the native PDF interface is given in Table 1.

When dealing with PDF uncertainties, the `nset` variable is used to distinguish the PDF grids corresponding to different eigenvectors. When CTEQ6.6M partons are used, `nset` varies in the range `nset=0,44`. When NNPDF2.0_100 partons are used, `nset` varies in the range `nset=0,100`. When MSTW2008 partons are used, `nset` varies in the range `nset=0,40` and the uncertainties we consider are those at 68% CL. When GJR08VF NLO or JR09VF NNLO partons are used, the `nset` variable should be in the range `nset=-13,13`. For A06 (ABKM09) NNLO partons, the variable `nset` should vary in the range `nset=0,23` (`nset=0,25`). The default choice is `nset=0`, corresponding to the central set. The variable `nset` is dummy when other PDF sets are used.

7 From Version 1.3 to 2.0

The main change in version 2.0 of the program is that the exact dependence on the masses of the top and bottom quarks up to NLO has been implemented. This is done by using the exact two-loop matrix element for $gg \rightarrow H$ [6], and the exact one loop

amplitudes for the real emission [7]. More details on the implementation of mass effects can be found in Ref. [3].

References

- [1] S. Catani and M. Grazzini, Phys. Rev. Lett. **98** (2007) 222002.
- [2] M. Grazzini, JHEP **0802** (2008) 043.
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- [4] S. Catani and M. H. Seymour, Phys. Lett. B **378** (1996) 287, Nucl. Phys. B **485** (1997) 291 [Erratum-ibid. B **510** (1998) 503].
- [5] J. Campbell, R.K. Ellis, *MCFM - Monte Carlo for FeMtobarn processes*, <http://mcfm.fnal.gov>
- [6] M. Spira, A. Djouadi, D. Graudenz and P. M. Zerwas, Nucl. Phys. B **453** (1995) 17 [hep-ph/9504378].
- [7] R. K. Ellis, I. Hinchliffe, M. Soldate and J. J. van der Bij, Nucl. Phys. B **297** (1988) 221.

iset	Pdf set	$\alpha_S(M_Z)$
1	CTEQ4 LO	0.132
2	CTEQ4 Standard NLO	0.116
11	MRST98 NLO central gluon	0.1175
12	MRST98 NLO higher gluon	0.1175
13	MRST98 NLO lower gluon	0.1175
14	MRST98 NLO lower α_S	0.1125
15	MRST98 NLO higher α_S	0.1225
16	MRST98 LO	0.125
21	CTEQ5M NLO Standard Msbar	0.118
22	CTEQ5D NLO DIS	0.118
23	CTEQ5L LO	0.127
24	CTEQ5HJ NLO Large-x gluon enhanced	0.118
25	CTEQ5HQ NLO Heavy Quark	0.118
28	CTEQ5M1 NLO Improved	0.118
29	CTEQ5HQ1 NLO Improved	0.118
30	MRST99 NLO	0.1175
31	MRST99 higher gluon	0.1175
32	MRST99 lower gluon	0.1175
33	MRST99 lower α_S	0.1125
34	MRST99 higher α_S	0.1225
41	MRST2001 NLO central gluon	0.119
42	MRST2001 NLO lower α_S	0.117
43	MRST2001 NLO higher α_S	0.121
44	MRST2001 NLO better fit to jet data	0.121
45	MRST2001 NNLO	0.1155
46	MRST2001 NNLO fast evolution	0.1155
47	MRST2001 NNLO slow evolution	0.1155
48	MRST2001 NNLO better fit to jet data	0.1180
51	CTEQ6L LO	0.118
52	CTEQ6L1 LO	0.130
53	CTEQ6M NLO	0.118
49	MRST2002 LO	0.130
61	MRST2002 NLO	0.1197
62	MRST2002 NNLO	0.1154
71	MRST2004 NLO	0.1205
72	MRST2004 NNLO	0.1167
90	MSTW2008 LO	0.13939
91	MSTW2008 NLO	0.12018
92	MSTW2008 NNLO	0.11707

Table 1: Available pdf sets and their corresponding `iset` and values of $\alpha_S(M_Z)$.